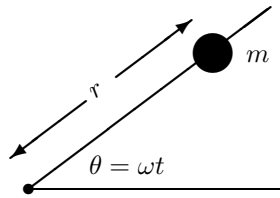


1 Bead on a Spinning Wire

This walks you through the problem.



A bead, with mass m , slides without friction on a long straight wire. The wire is spinning about one end at a constant angular speed ω . Let r be the position of the bead on the wire as measured from the point of rotation of the wire. There is no gravity. All the motion is in a plane.

1.1 Lagrangian

In terms of r , \dot{r} , m , ω , and t , write an expression for the Lagrangian $L(r, \dot{r}, t)$ for this system. Note that $U = 0$.

1.1 solution

$$L = T - U = \frac{1}{2}m(\dot{r}^2 + r^2\omega^2) - 0 = \frac{1}{2}m\dot{r}^2 + \frac{1}{2}mr^2\omega^2 \Rightarrow L(r, \dot{r}) = \frac{1}{2}m\dot{r}^2 + \frac{1}{2}mr^2\omega^2. \quad (1.1)$$

1.2 Lagrange's Equation of Motion

Using your Lagrangian, find the equation of motion for r .

1.2 solution

$$\frac{\partial L}{\partial r} - \frac{d}{dt} \frac{\partial L}{\partial \dot{r}} = 0 \Rightarrow mr\omega^2 - \frac{d}{dt}m\dot{r} = 0 \Rightarrow \ddot{r} = r\omega^2. \quad (1.2)$$

1.3 Canonical Conjugate Momentum

In general, the canonical momentum that is conjugate (paired) to generalized position q_i is $p_{q_i} \equiv p_i = \frac{\partial L}{\partial \dot{q}_i}$.

The canonical momentum that is conjugate to r is $p_r \equiv \frac{\partial L}{\partial \dot{r}}$. Find the canonical momentum that is conjugate to r , p_r , as a function of r , \dot{r} , ω , and t .

1.3 solution

$$p_r \equiv \frac{\partial L}{\partial \dot{r}} = \frac{\partial}{\partial \dot{r}} \left(\frac{1}{2}m\dot{r}^2 \right) \Rightarrow p_r = m\dot{r}. \quad (1.3)$$

1.4 Hamiltonian

The Hamiltonian in general can be defined as

$$H(q_k, p_k, t) = \sum_j p_j \dot{q}_j - L(q_k, \dot{q}_k, t), \tag{1.4}$$

where the Hamiltonian must be written as a function of just the dependent dynamical variables q_k , p_k , and t . There must be no \dot{q}_k dependence in H . So we are assuming that \dot{q}_k can be written as a function of q_k , p_k , and t , in order to make this transformation.

In our case $H(r, p_r) = p_r \dot{r} - L(r, \dot{r})$. In terms of r , p_r , m , and ω , find the Hamiltonian, $H(r, p_r, t)$ for this system.

1.4 solution

From this definition of H we get

$$H = p_r \dot{r} - \left(\frac{1}{2} m \dot{r}^2 + \frac{1}{2} m r^2 \omega^2 \right) = p_r \left(\frac{p_r}{m} \right) - \frac{1}{2} m \left(\frac{p_r}{m} \right)^2 - \frac{1}{2} m r^2 \omega^2 \Rightarrow H(r, p_r) = \frac{p_r^2}{2m} - \frac{1}{2} m r^2 \omega^2. \tag{1.5}$$

1.5 Hamilton's Equations of Motion

In Lagrangian dynamics, for each generalized position, q_i , ($i = 1, 2, \dots, s$) there is an equation of motion that is a second-order ordinary differential equation. In Hamiltonian Dynamics, for each corresponding second-order differential equation of Lagrangian dynamics there are two first-order ordinary differential equations, one of the generalized position, q_i , and one for the corresponding canonical momentum, $p_{q_i} \equiv p_i$. These Hamiltonian equations of motion are gotten from the Hamiltonian like so

$$\dot{q}_i = \frac{\partial H}{\partial p_{q_i}} \quad \text{and} \quad \dot{p}_i = -\frac{\partial H}{\partial q_i} \tag{1.6}$$

Use the Hamiltonian that you just calculated and equations 1.6 to find the equations of motion for r and p_r , and show that these equations of motion for r and p_r are equivalent to the equation of motion from section 1.2.

1.5 solution

$$\dot{r} = \frac{\partial H}{\partial p_r} = \frac{p_r}{m} \Rightarrow \dot{r} = \frac{p_r}{m}, \tag{1.7}$$

and

$$\dot{p}_r = -\frac{\partial H}{\partial r} = -(-mr\omega^2) \Rightarrow \dot{p}_r = mr\omega^2. \tag{1.8}$$

From equation 1.7 and equation 1.8

$$\ddot{r} = \frac{1}{m} \dot{p}_r = \frac{1}{m} (mr\omega^2) = r\omega^2 \Rightarrow \ddot{r} = r\omega^2, \tag{1.9}$$

which is equation 1.2. So the equations of motion from Lagrangian dynamics and Hamiltonian Dynamics are equivalent.

1.6 Energy vs Hamiltonian

In terms of r , p_r , m , ω , and t , **(a)** find the difference between the total energy E and the Hamiltonian H for this system ($E - H$). **(b)** Is H constant in time? **(c)** Is E constant in time? Note that $E = T = L$. Explain your results.

1.6 solution

$$E - H = \left(\frac{1}{2} m \dot{r}^2 + \frac{1}{2} m r^2 \omega^2 \right) - \left(\frac{p_r^2}{2m} - \frac{1}{2} m r^2 \omega^2 \right) = \frac{1}{2} m \dot{r}^2 + \frac{1}{2} m r^2 \omega^2 - \frac{1}{2} m \dot{r}^2 + \frac{1}{2} m r^2 \omega^2 = m r^2 \omega^2$$

$$\Rightarrow \boxed{\text{(a) } E - H = m r^2 \omega^2}.$$
 (1.10)

$$\frac{dH}{dt} = \frac{d}{dt} \left(\frac{p_r^2}{2m} - \frac{1}{2} m r^2 \omega^2 \right) = \frac{p_r}{m} \dot{p}_r - m \omega^2 r \dot{r} = \frac{p_r}{m} (m r \omega^2) - m \omega^2 r \left(\frac{p_r}{m} \right) = 0,$$
 (1.11)

where we have used the equations of motion for $\dot{p}_r = m r \omega^2$ and $\dot{r} = \frac{p_r}{m}$. So $\boxed{\text{(b) } H \text{ is constant}}$.

$$\frac{dE}{dt} = \frac{d}{dt} \left(\frac{1}{2} m \dot{r}^2 + \frac{1}{2} m r^2 \omega^2 \right) = \frac{p_r}{m} \dot{p}_r + m \omega^2 r \dot{r} = 2 r p_r \omega^2,$$
 (1.12)

So $\boxed{\text{(c) } E \text{ is not constant}}$.