

1 $r(\theta)$ – Orbital Path for Another Force

A particle of mass m moves with angular momentum l and with total energy E about a fixed center with a force

$$F(r) = -\frac{k}{r^2} + \frac{\lambda}{r^3} \quad (1.1)$$

where k and λ are greater than zero, and r is the distance from the particle to the center. **(a)** Show that the equation for the orbit, $r(\theta)$, may have the form

$$\frac{\alpha}{r} = 1 + \epsilon \cos(\beta\theta), \quad (1.2)$$

finding the constants α , ϵ , and β in terms of the given quantities m , k , λ , l , and E . Assume that the potential energies are zero at $r = \infty$ in defining the total energy, E . **(b)** For what values of β is the orbit closed?

1.0 solution

(a) I'll show all the detail here, but you will likely only need a subset of this by using the results of the Kepler problem or other short-cuts.

$$d\theta = \frac{d\theta}{dt} \frac{dt}{dr} dr = \frac{\dot{\theta}}{\dot{r}} dr. \quad (1.3)$$

$$U(r) = -\int \left(-\frac{k}{r^2} + \frac{\lambda}{r^3} \right) dr = -\frac{k}{r} + \frac{\lambda}{2r^2}, \quad (1.4)$$

where we set $U(\infty) = 0$. From conservation of energy and angular momentum, $l = mr^2\dot{\theta}$,

$$\begin{aligned} E &= \frac{1}{2}m\dot{r}^2 + \frac{1}{2}mr^2\dot{\theta}^2 - \frac{k}{r} + \frac{\lambda}{2r^2} = \frac{1}{2}m\dot{r}^2 + \frac{1}{2}mr^2 \left(\frac{l^2}{m^2r^4} \right) - \frac{k}{r} + \frac{\lambda}{2r^2} = \frac{1}{2}m\dot{r}^2 + \frac{l^2}{2mr^2} - \frac{k}{r} + \frac{\lambda}{2r^2} \\ \Rightarrow E &= \frac{1}{2}m\dot{r}^2 + \left(\frac{l^2}{2m} + \frac{\lambda}{2} \right) \frac{1}{r^2} - \frac{k}{r} \quad \Rightarrow \quad \dot{r} = \pm \sqrt{\frac{2}{m} \left[E - \left(\frac{l^2}{2m} + \frac{\lambda}{2} \right) \frac{1}{r^2} - \frac{k}{r} \right]} \end{aligned} \quad (1.5)$$

This with equation 1.3 gives

$$\Rightarrow d\theta = \pm \frac{\frac{l}{mr^2} dr}{\sqrt{\frac{2}{m} \left[E - \left(\frac{l^2}{2m} + \frac{\lambda}{2} \right) \frac{1}{r^2} + \frac{k}{r} \right]}} = \pm \frac{l}{\sqrt{2m}} \frac{\frac{dr}{r^2}}{\sqrt{E - \left(\frac{l^2}{2m} + \frac{\lambda}{2} \right) \frac{1}{r^2} + \frac{k}{r}}}. \quad (1.6)$$

With the change of variables $u = \frac{1}{r}$, $\frac{dr}{r^2} = -du$, we get

$$d\theta = \mp \frac{l}{\sqrt{2m}} \frac{du}{\sqrt{E + ku - \left(\frac{l^2}{2m} + \frac{\lambda}{2} \right) u^2}}. \quad (1.7)$$

We complete the square of the term

$$\left(\frac{l^2}{2m} + \frac{\lambda}{2} \right) u^2 - ku,$$

giving

$$d\theta = \mp \frac{l}{\sqrt{2m}} \frac{du}{\sqrt{\left(E + \frac{k^2}{2\frac{l^2}{m} + 2\lambda} \right) - \left(\sqrt{\frac{l^2}{2m} + \frac{\lambda}{2}} u - \frac{k}{\sqrt{2\frac{l^2}{m} + 2\lambda}} \right)^2}}. \quad (1.8)$$

Change variables to

$$y = \sqrt{\frac{l^2}{2m} + \frac{\lambda}{2}} u - \frac{k}{\sqrt{\frac{2l^2}{m} + 2\lambda}} \quad \text{so} \quad dy = \sqrt{\frac{l^2}{2m} + \frac{\lambda}{2}} du \quad (1.9)$$

$$\Rightarrow \theta = \mp \frac{1}{\sqrt{\frac{l^2}{2m} + \frac{\lambda}{2}}} \frac{l}{\sqrt{2m}} \int \frac{dy}{\sqrt{\left[\left(E + \frac{k^2}{2\frac{l^2}{m} + 2\lambda} \right) - y^2 \right]}} = \mp \frac{1}{\sqrt{\frac{l^2}{2m} + \frac{\lambda}{2}}} \frac{l}{\sqrt{2m}} \cos^{-1} \left(\frac{y}{\sqrt{E + \frac{k^2}{2\frac{l^2}{m} + 2\lambda}}} \right), \quad (1.10)$$

where we have set the integration constant to zero, so that $\theta = 0$ will have a minimum r value.

$$\begin{aligned} \Rightarrow \cos \left(\sqrt{\frac{1}{2m} + \frac{\lambda}{2l^2}} \sqrt{2m} \theta \right) &= \frac{\sqrt{\frac{l^2}{2m} + \frac{\lambda}{2}} u - \frac{k}{\sqrt{\frac{2l^2}{m} + 2\lambda}}}{\sqrt{E + \frac{k^2}{2\frac{l^2}{m} + 2\lambda}}} \\ \Rightarrow \frac{k}{\sqrt{\frac{2l^2}{m} + 2\lambda}} + \sqrt{E + \frac{k^2}{2\frac{l^2}{m} + 2\lambda}} \cos \left(\sqrt{1 + \frac{m\lambda}{l^2}} \theta \right) &= \sqrt{\frac{l^2}{2m} + \frac{\lambda}{2}} u, \\ \Rightarrow \frac{\frac{l^2 + m\lambda}{mk}}{r} = 1 + \frac{\sqrt{\frac{2l^2}{m} + 2\lambda}}{k} \sqrt{E + \frac{k^2}{2\frac{l^2}{m} + 2\lambda}} \cos \left(\sqrt{1 + \frac{m\lambda}{l^2}} \theta \right) \\ \Rightarrow \boxed{\frac{\frac{l^2 + m\lambda}{mk}}{r} = 1 + \left[\sqrt{\frac{2E(l^2 + m\lambda)}{k^2} + 1} \right] \cos \left(\sqrt{1 + \frac{m\lambda}{l^2}} \theta \right)}. \end{aligned} \quad (1.11)$$

So

$$\boxed{\alpha \equiv \frac{l^2 + m\lambda}{mk}}, \quad \boxed{\beta \equiv \sqrt{1 + \frac{m\lambda}{l^2}}} \quad \text{and} \quad \boxed{\epsilon \equiv \sqrt{\frac{2E(l^2 + m\lambda)}{k^2} + 1}}, \quad (1.12)$$

and

$$\frac{\alpha}{r} = 1 + \epsilon \cos(\beta\theta) \quad (1.13)$$

(b) In order for there to be a closed orbit the phase in cosine term, $\beta\theta$, must be $n2\pi$ when $\theta = m2\pi$ where m and n are integers, and the orbit will close after m revolutions. Therefore, in order for there to be a closed orbit

$$\beta(m2\pi) = n2\pi \quad \Rightarrow \quad \beta = \frac{n}{m} = \text{rational number} \quad \Rightarrow \quad \sqrt{1 + \frac{m\lambda}{l^2}} = \text{rational number}. \quad (1.14)$$

